

Ekurhuleni North	District Paper – Mathematics Paper 2 Grade 11 Marking Guideline	June 2023
<b>SOLUTION</b> <b>NB:</b> Accept valid alternate solutions. <b>Consistent Accuracy</b> applies to the whole memo.		
<b>QUESTION 1</b>		
1.1	$m_{HM} = \frac{6 - 4}{3 + 1}$ $= \frac{1}{2}$	✓ subs into formula ✓ answer
1.2	$R(x; y) = \left( \frac{3 + 4}{2}; \frac{6 + 1}{2} \right)$ $= \left( \frac{7}{2}; \frac{7}{2} \right)$	✓ Correct Substitution  ✓ Answer
1.3	$m_{HM} = m_{AP}$ $\tan \theta = \frac{1}{2}$ $\theta = 26,57^\circ$	✓ $HM \parallel AT$ ✓ Correct substitution ✓ Answer
1.4	<p><math>R</math> is the midpoint of diagonal <math>AH</math> (Diags of a parm)          Let the coordinates of <math>A</math> be <math>(x; y)</math>.</p> $\frac{-1 + x}{2} = \frac{7}{2}$ $x = 8$ $\frac{4 + y}{2} = \frac{7}{2}$ $y = 3$ <p><math>\therefore A(8; 3)</math></p>	$\checkmark \frac{-1+x}{2} = \frac{7}{2}$ $\checkmark 8$ $\checkmark \frac{4+y}{2} = \frac{7}{2}$ $\checkmark 3$
1.5	$m_{AT} = \frac{1}{2} \text{ (parallel lines)}$ $y = \frac{1}{2}x + c$ <p>Sub <math>T(4; 1)</math></p>	    ✓ Sub $T(4; 1)$   ✓ $c = -1$

# Maths Paper1 2013 Memo

**David Dos Santos Ferreira, Wolfgang  
Staubach**



## Maths Paper1 2013 Memo:

**A Complete Classification of the Isolated Singularities for Nonlinear Elliptic Equations with Inverse Square Potentials** Florica C. Cîrstea, 2014-01-08 In particular for  $b \neq 0$  we find a sharp condition on  $h$  such that the origin is a removable singularity for all non negative solutions of  $\Delta u = h(x)|u|^{p-2}u$  thus addressing an open question of Vazquez and Vron

**Singularity Theory for Non-Twist KAM Tori** A. González-Enríquez, A. Haro, R. de la Llave, 2014-01-08 In this monograph the authors introduce a new method to study bifurcations of KAM tori with fixed Diophantine frequency in parameter dependent Hamiltonian systems It is based on Singularity Theory of critical points of a real valued function which the authors call the potential The potential is constructed in such a way that nondegenerate critical points of the potential correspond to twist invariant tori i.e. with nondegenerate torsion and degenerate critical points of the potential correspond to non twist invariant tori Hence bifurcating points correspond to non twist tori

**Global and Local Regularity of Fourier Integral Operators on Weighted and Unweighted Spaces** David Dos Santos Ferreira, Wolfgang Staubach, 2014-04-07 The authors investigate the global continuity on spaces with of Fourier integral operators with smooth and rough amplitudes and or phase functions subject to certain necessary non degeneracy conditions In this context they prove the optimal global boundedness result for Fourier integral operators with non degenerate phase functions and the most general smooth Hörmander class amplitudes i.e. those in  $S_{\rho, \delta}^{-\infty}$  with  $\rho \geq \delta$  They also prove the very first results concerning the continuity of smooth and rough Fourier integral operators on weighted spaces with and i.e. the Muckenhoupt weights for operators with rough and smooth amplitudes and phase functions satisfying a suitable rank condition

**Large Deviations for Additive Functionals of Markov Chains** Alejandro D. de Acosta, Peter Ney, 2014-03-05

**Combinatorial Floer Homology** Vin de Silva, Joel W. Robbin, Dietmar A. Salamon, 2014-06-05 The authors define combinatorial Floer homology of a transverse pair of noncontractible nonisotopic embedded loops in an oriented manifold without boundary prove that it is invariant under isotopy and prove that it is isomorphic to the original Lagrangian Floer homology Their proof uses a formula for the Viterbo Maslov index for a smooth lune in a manifold

**Weighted Bergman Spaces Induced by Rapidly Increasing Weights** Jose Angel Pelaez, Jouni Rattya, 2014-01-08 This monograph is devoted to the study of the weighted Bergman space  $A_{p, \omega}$  of the unit disc  $\mathbb{D}$  that is induced by a radial continuous weight  $\omega$  satisfying  $\lim_{r \rightarrow 1} \frac{1}{2\pi} \int_0^{2\pi} \omega(re^{i\theta}) d\theta = 1$  Every such  $A_{p, \omega}$  lies between the Hardy space  $H^p$  and every classical weighted Bergman space  $A_{p, \alpha}$  Even if it is well known that  $H^p$  is the limit of  $A_{p, \alpha}$  as  $\alpha \rightarrow 1$  in many respects it is shown that  $A_{p, \omega}$  lies closer to  $H^p$  than any  $A_{p, \alpha}$  and that several finer function theoretic properties of  $A_{p, \alpha}$  do not carry over to  $A_{p, \omega}$

**Nonlinear Stability of Ekman Boundary Layers in Rotating Stratified Fluids** Hajime Koba, 2014-03-05 A stationary solution of the rotating Navier Stokes equations with a boundary condition is called an Ekman boundary layer This book constructs stationary solutions of the rotating Navier Stokes Boussinesq equations with stratification effects in the case when the

rotating axis is not necessarily perpendicular to the horizon The author calls such stationary solutions Ekman layers This book shows the existence of a weak solution to an Ekman perturbed system which satisfies the strong energy inequality Moreover the author discusses the uniqueness of weak solutions and computes the decay rate of weak solutions with respect to time under some assumptions on the Ekman layers and the physical parameters The author also shows that there exists a unique global in time strong solution of the perturbed system when the initial datum is sufficiently small Comparing a weak solution satisfying the strong energy inequality with the strong solution implies that the weak solution is smooth with respect to time when time is sufficiently large

Near Soliton Evolution for Equivariant Schrodinger Maps in Two Spatial Dimensions Ioan Bejenaru, Daniel Tataru, 2014-03-05 The authors consider the Schrödinger Map equation in 2+1 dimensions with values into  $S^2$  This admits a lowest energy steady state  $Q$  namely the stereographic projection which extends to a two dimensional family of steady states by scaling and rotation The authors prove that  $Q$  is unstable in the energy space  $\dot{H}^1$  However in the process of proving this they also show that within the equivariant class  $Q$  is stable in a stronger topology  $X \subset \dot{H}^1$

*Advances and Applications of DSMT for Information Fusion, Vol. IV* Florentin Smarandache, Jean Dezert, 2015-03-01 The fourth volume on Advances and Applications of Dezert Smarandache Theory DSMT for information fusion collects theoretical and applied contributions of researchers working in different fields of applications and in mathematics The contributions see List of Articles published in this book at the end of the volume have been published or presented after disseminating the third volume 2009 <http://fs.gallup.unm.edu/DSMT/book3.pdf> in international conferences seminars workshops and journals

**Advances and Applications of DSMT for Information Fusion. Collected Works, Volume 4** Florentin Smarandache, Jean Dezert, 2015-07-01 The fourth volume on Advances and Applications of Dezert Smarandache Theory DSMT for information fusion collects theoretical and applied contributions of researchers working in different fields of applications and in mathematics The contributions have been published or presented after disseminating the third volume 2009 <http://fs.gallup.unm.edu/DSMT/book3.pdf> in international conferences seminars workshops and journals

This book delves into Maths Paper1 2013 Memo. Maths Paper1 2013 Memo is a vital topic that must be grasped by everyone, from students and scholars to the general public. The book will furnish comprehensive and in-depth insights into Maths Paper1 2013 Memo, encompassing both the fundamentals and more intricate discussions.

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- Chapter 3: Maths Paper1 2013 Memo in Everyday Life
- Chapter 4: Maths Paper1 2013 Memo in Specific Contexts
- Chapter 5: Conclusion

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